

“HIGH ENERGY ASTROPHYSICS” (A subjective introduction)

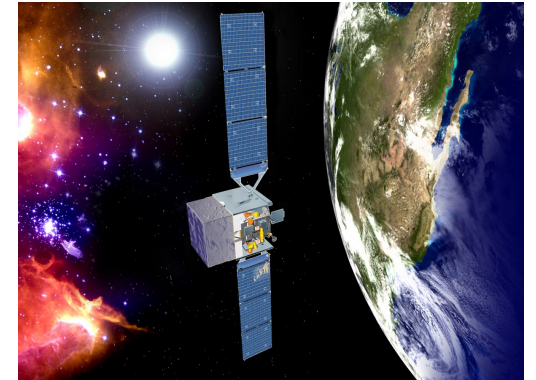
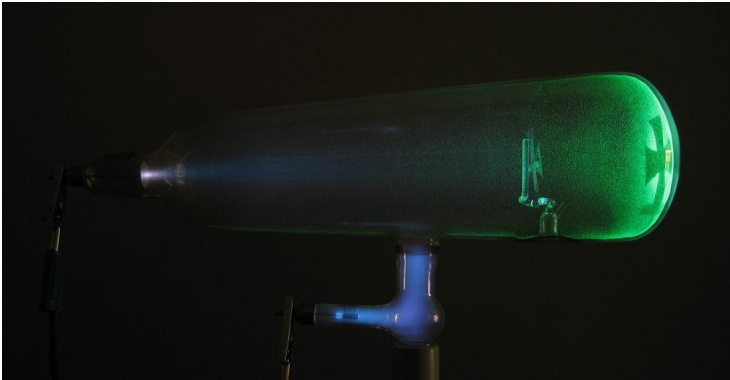
II. Special relativity - reminder

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Story so far ...

High energy astrophysics: “ ... areas of research undertaken by astrophysicists, physicists and astronomers interested in high energy phenomena of all types ”

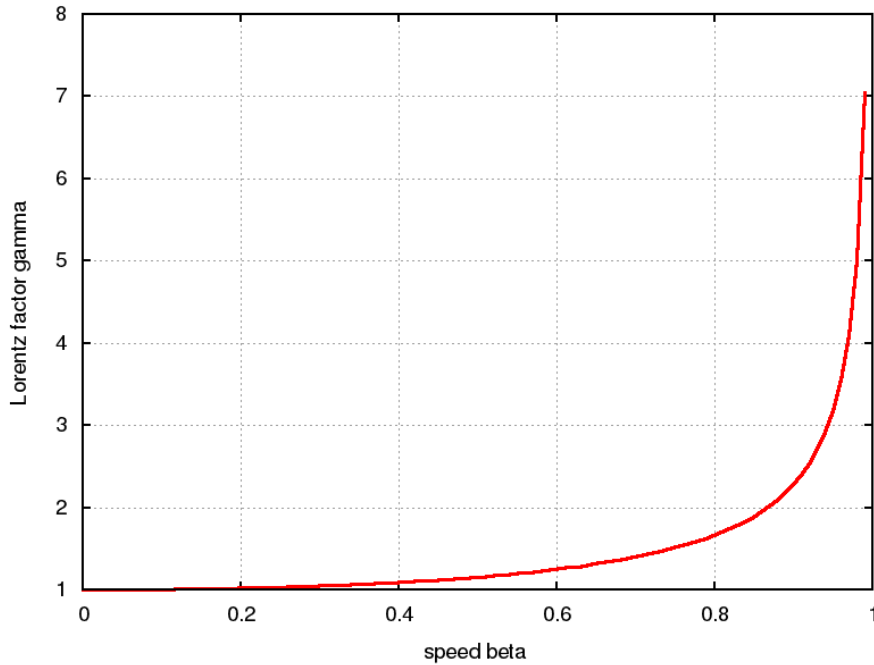


gamma-ray band: $\nu \gg 3 \times 10^{19}$ Hz; $E \gg 100$ keV

Speed, energy, Lorentz factor

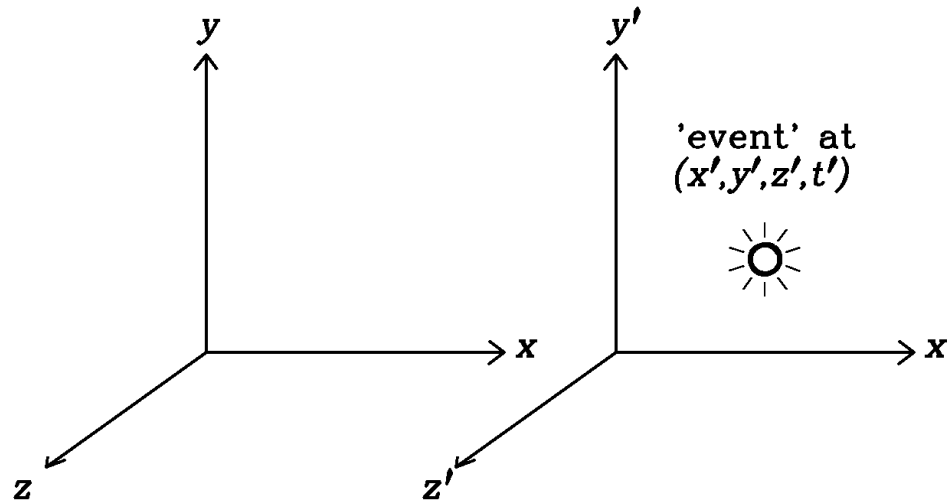
$$\beta = \frac{v}{c}, \quad \text{Lorentz factor } \gamma = \sqrt{\frac{1}{1-\beta^2}}$$

$\gamma = 1$ for $\beta = 0$ $\gamma = \infty$ for $\beta \rightarrow 1$



β	γ
0.87	2
0.995	10
0.9999995	1000
...	10^7
...	10^{11}

Lorentz transformation



space-time event in K – (ct, x, y, z) , in K' – (ct', x', y', z')

Lorentz transformation

$$ct' = \gamma(ct - \beta x)$$
$$x' = \gamma(x - \beta ct), \quad y' = y, \quad z' = z$$

inverse Lorentz transformation

$$ct = \gamma(ct' + \beta x')$$
$$x = \gamma(x' + \beta ct'), \quad y = y', \quad z = z'$$

Length contraction and time dilation

- *length contraction*

$$L' = x'_2 - x'_1 = \gamma (x_2 - x_1) = \gamma L$$
$$t'_2 = t'_1 \quad \text{and} \quad t_2 = t_1$$

- *time dilation*

$$T = t_2 - t_1 = \gamma (t'_2 - t'_1) = \gamma T'$$
$$x'_2 = x'_1$$

Proper time

proper time τ

$$c^2 d\tau^2 = c^2 dt^2 - (dx^2 + dy^2 + dz^2)$$

$$d\tau' = ?$$

proper time is a *Lorentz invariant*

Transformation of velocities

for velocity components u'_x, u'_y, u'_z

$$u_x = \frac{dx}{dt} = \frac{\gamma(dx' + \beta c dt')}{\gamma(dt' + \beta dx'/c)} = \frac{dx'/dt' + \beta c}{1 + \beta/c dx'/dt'} = \frac{u'_x + \beta c}{1 + \beta/c u'_x}$$

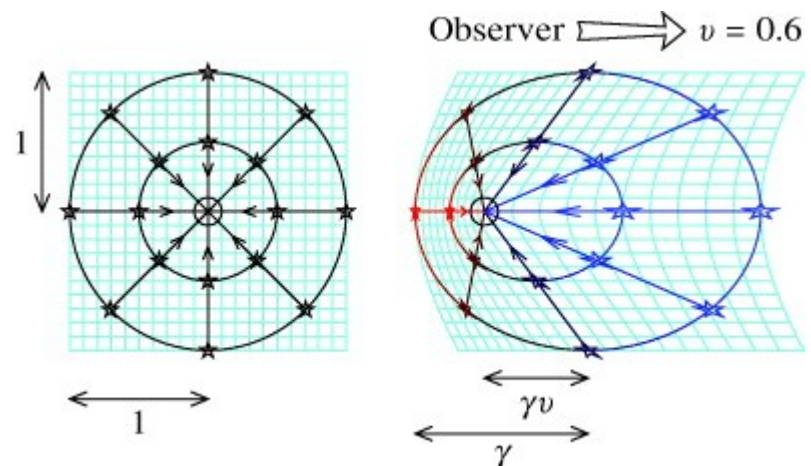
$$u_y = \frac{u'_y}{\gamma(1 + \beta/c u'_x)}, \quad u_z = \frac{u'_z}{\gamma(1 + \beta/c u'_x)}$$

for perpendicular and parallel components

$$u_{\parallel} = \frac{u'_{\parallel} + \beta c}{1 + \beta/c u'_{\parallel}}, \quad u_{\perp} = \frac{u_{\perp}}{1 + \beta/c u'_{\parallel}}$$

relativistic aberration

$$\tan \theta = \frac{u_{\perp}}{u_{\parallel}} = \frac{u' \sin \theta'}{\gamma(u' \cos \theta' + \beta c)}$$



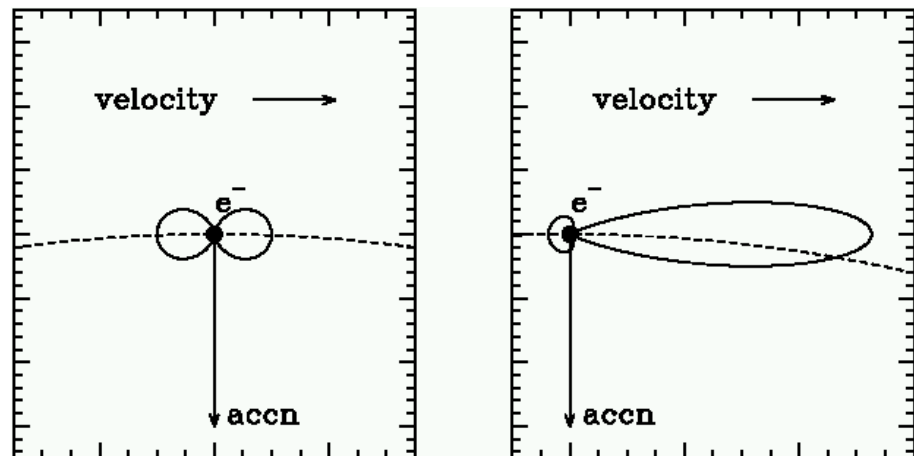
Relativistic beaming

Consider a photon ($u'=c$) emitted perpendicularly ($\Theta'=\pi/2$) to the direction of motion:

$$\tan \theta = \frac{c \sin \pi/2}{\gamma (c \cos \pi/2 + \beta c)} = \frac{1}{\beta \gamma}$$

for large velocities

$$\theta \approx \tan \theta \approx \frac{1}{\gamma}$$



Doppler effect

non-relativistic case $\omega_{obs} = \frac{1}{(1 - \beta \cos \theta)} \omega_{source}$

relativistic case

$$\omega_{source} = \omega' = \frac{2\pi}{\Delta T'}$$

$$\Delta T = \gamma \Delta T' = \gamma \frac{2\pi}{\omega'}$$

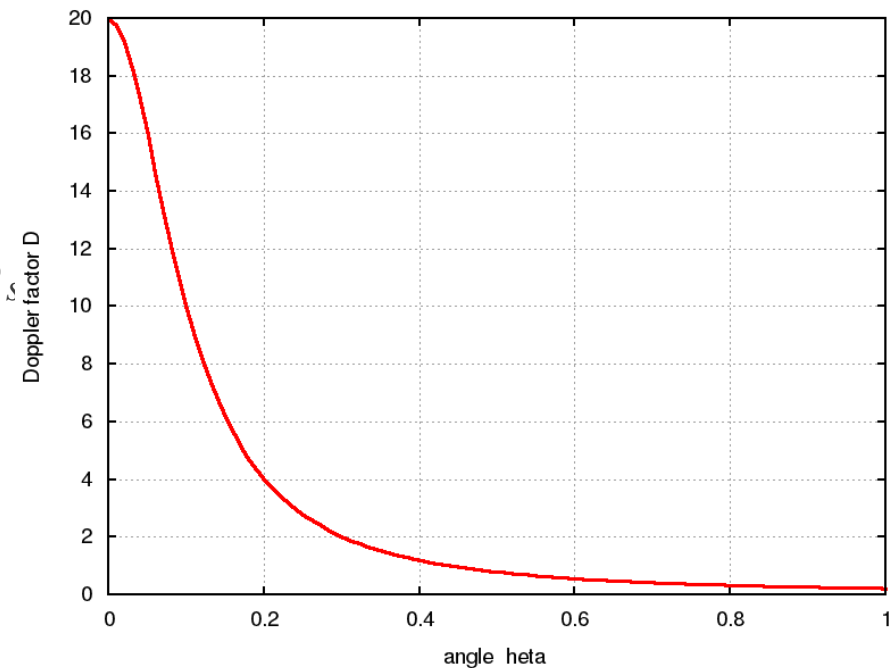
$$\Delta T_{obs} = \Delta T - \frac{d}{c} = \Delta T (1 - \beta \cos \theta)$$

$$\omega_{obs} = \frac{2\pi}{\Delta T_{obs}} = \frac{1}{\gamma (1 - \beta \cos \theta)} \omega_{source} = D \omega_{source}$$

Doppler factor D

$$D = \frac{1}{\gamma (1 - \beta \cos \theta)}$$

$$\theta = 0 \Rightarrow D = 2\gamma, \theta = 1/\gamma \Rightarrow D = \gamma, \theta = \pi/2 \Rightarrow D = 1/\gamma$$



Tensor calculus

vector

$$\vec{x} = (x^1, x^2, x^3) = x^1 \hat{e}_1 + x^2 \hat{e}_2 + x^3 \hat{e}_3 = x^i \hat{e}_i$$

Einstein summation convention

$$a_i b^i = \sum_{i=1}^{i=3} a_i b^i$$

scalar product

$$\vec{x} \circ \vec{y} = (x^i \hat{e}_i) \circ (y^k \hat{e}_k) = (\hat{e}_i \circ \hat{e}_j) x^i y^j = g_{ij} x^i y^j$$

metric tensor

$$g = (g_{ij}) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Tensor calculus

four-vectors

$$X^\mu = (x^0, x^1, x^2, x^3) = x^0 \hat{e}_0 + x^1 \hat{e}_1 + x^2 \hat{e}_2 + x^3 \hat{e}_3 = x^\mu \hat{e}_\mu$$

convention

$$i, j, k, \dots \text{run } 1, 2, 3 \quad \mu, \nu, \xi, \dots \text{run } 0, 1, 2, 3$$

special relativity metric tensor

$$\eta = (\eta_{ij}) = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

scalar product

$$c^2 d\tau = -\eta_{\mu\nu} dx^\mu dx^\nu \equiv -ds^2$$

Tensor calculus

contra- and covariant coordinates

$$X_{\alpha} = g_{\alpha\beta} X^{\beta}, \quad X^{\alpha} = g^{\alpha\beta} X_{\beta}, \quad g^{\alpha\lambda} g_{\lambda\beta} = \delta_{\beta}^{\alpha}$$

contravariant vector transformation

$$x'^{\mu} = \Lambda_{\nu}^{\mu} x^{\nu}, \quad \Lambda_{\nu}^{\mu} = \frac{\partial x'^{\mu}}{\partial x^{\nu}}$$

Lorentz transformation

$$(x'^{\mu}) = \begin{pmatrix} ct' \\ x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \gamma & -\gamma\beta & 0 & 0 \\ -\gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix}$$

$$dx'^{\mu} = \Lambda_{\nu}^{\mu} dx^{\nu}$$

covariant vector transformation

$$x'_{\mu} = \tilde{\Lambda}_{\mu}^{\nu} x_{\nu}, \quad \tilde{\Lambda}_{\mu}^{\xi} \Lambda_{\xi}^{\nu} = \delta_{\mu}^{\nu}$$

Tensor calculus

Lorentz invariance

$$A'_{\mu} B'^{\mu} = (\tilde{\Lambda}_{\mu}^{\xi} A_{\xi}) (\Lambda_{\lambda}^{\mu} B^{\lambda}) = \tilde{\Lambda}_{\mu}^{\xi} \Lambda_{\lambda}^{\mu} A_{\xi} B^{\lambda} = \delta_{\lambda}^{\xi} A_{\xi} B^{\lambda} = A_{\lambda} B^{\lambda} = A_{\mu} B^{\mu}$$

Example

$$dx^{\mu} dx_{\mu} = \eta_{\mu\nu} dx^{\mu} dx^{\nu} = -c^2 d\tau = ds^2$$

Four-vectors

four-velocity

$$(U^\mu) \equiv \frac{dx^\mu}{d\tau} = \gamma \frac{dx^\mu}{dt} = \gamma \left(\frac{dct}{dt}, \frac{dx}{dt}, \frac{dy}{dt}, \frac{dz}{dt} \right) = \gamma (c, \vec{v})$$

$$U_\mu U^\mu = \frac{dx_\mu}{d\tau} \frac{dx^\mu}{d\tau} = \frac{-c^2 d\tau^2}{d\tau^2} = -c^2$$

four-momentum

$$(p^\mu) \equiv m_0 (U^\mu) = (\gamma m_0 c, \gamma m_0 \vec{v}) = \left(\frac{\gamma m_0 c^2}{c}, \gamma m_0 \vec{v} \right) = \left(\frac{E}{c}, \vec{p} \right)$$

$$p_\mu p^\mu = \eta_{\mu\nu} p^\nu p^\mu = \frac{-E^2}{c^2} + p^2 = -m_0^2 c^2$$

for photon

$$\vec{p} = \frac{E}{c} \hat{e}_p, \quad p_\mu p^\mu = \frac{-E^2}{c^2} + p^2 = 0$$

four-acceleration and four-force

$$a^\mu = \frac{dU^\mu}{d\tau}, \quad F^\mu = m_0 a^\mu = \frac{dp^\mu}{d\tau}.$$

Relativistic electrodynamics

four-current and *four-potential*

$$(j^\mu) = (\rho c, \vec{j}), \quad (A^\mu) = (\Phi, \vec{A})$$

electromagnetic field tensor

$$F_{\mu\nu} \equiv A_{\nu,\mu} - A_{\mu,\nu} = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{pmatrix}$$

Maxwell equations

$$\partial^\mu F_{\nu\mu} = \frac{4\pi}{c} j^\nu$$

Exercises

1. Compton scattering

Consider a photon scattered off an electron. Use four-momentum conservation law to show that the shift in a photon wavelength $\Delta\lambda$ is

$$\Delta\lambda = \lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta) ,$$

where θ is a scattering angle.

2. Can an isolated electron absorb a photon?